

the distance $c \text{ } \mathbf{1}$ ($=C \text{ } \mathbf{1}$) from $O \text{ } P$, or in the direction $O \text{ } \not{p}$ parallel to $c \text{ } a$. In virtue of both movements, then, it will be found at x , the extremity of the diagonal of the square $O \text{ } x$ at that moment. And similarly at the end of the $2d$, $3d$, &c., interval, it will be found at the extremities of the diagonals of the squares next in succession, and as these all lie in one line, $O \text{ } E$, 45° inclined both to $O \text{ } P$ and $O \text{ } \not{p}$, it appears that in this case the resultant vibration will be rectilinear, and will be performed along the diagonal $E \text{ } G$ of the square $E \text{ } F \text{ } G \text{ } H$; and thus it appears that the superposition of two rays of equal intensity, polarized in opposite (*i.e.*, rectangular) planes, results in the production of a ray polarized in a plane 45° inclined to each of the former. Moreover, the square of the diagonal being double that of either side of a square, and the intensity of a ray being measured by *the square* of the vibrational excursion of its ethereal molecules, the intensity of the compound ray will be double that of the components, or, equal to their sum. And, *vice versâ*, any polarized ray may be considered as equivalent to two rays, each of half its intensity, polarized in planes 45° inclined on one side, and on the other of its plane of polarization. It need hardly be observed that if the molecule in starting from O be moving in the direction $C \text{ } A$, in virtue of the one vibration, and of $c \text{ } b$ in virtue of the other,—that is, if it be commencing its first semi-vibration in the one direction, and its second in the other, or again in other words, if the vibrations differ in phase by an exact semi-undulation; all the same reasoning will apply, with this